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1. We found in class that gravity waves satisfy

$$\phi = \frac{a\omega}{k} \frac{\cosh k(z+H)}{\sinh kH} \sin(kx - \omega t)$$

where ϕ is the "velocity potential", $u = \frac{\partial \phi}{\partial x}$, $w = \frac{\partial \phi}{\partial z}$

We're interested in the limit $kH \gg 1$. Expanding terms

in the relevant part of ϕ , we can write

$$f \equiv \frac{\cosh k(z+H)}{\sinh kH} = \frac{e^{kH(1+z/H)} + e^{-kH(1+z/H)}}{e^{kH} + e^{-kH}} \approx \frac{e^{kH(1+z/H)}}{e^{kH}} = e^{kz}$$

where I have taken the $kH \gg 1$ limit holding the fractional depth ($\frac{z}{H}$) constant; the approximation holds anywhere above the bottom ($z = -H$).

In this regime,

$$\phi \approx \frac{a\omega}{k} e^{kz} \sin(kx - \omega t)$$

so the velocities are

$$(u, w) = \left(\frac{\partial \phi}{\partial x}, \frac{\partial \phi}{\partial z} \right) = a\omega e^{kz} (\cos(kx - \omega t), \sin(kx - \omega t))$$

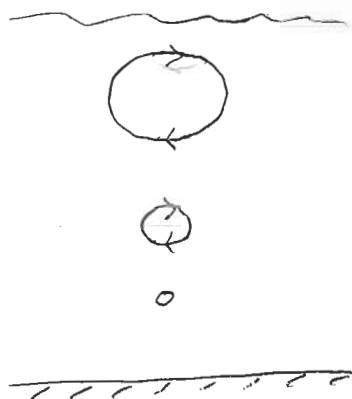
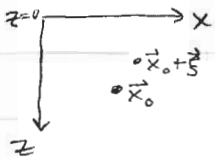
Consider a small amplitude trajectory \vec{s} around \vec{x}_0 ($k|\vec{s}| \ll 1$).

We can assume the velocity field is fairly constant along this trajectory, so

$$\vec{s} = \int \vec{u}(\vec{x}_0) \cdot dt = a e^{kz_0} (-\sin(kx_0 - \omega t), \cos(kx_0 - \omega t))$$

Which is a circular trajectory with

radius $a e^{kz_0} = a e^{-k|z_0|}$ (since $z_0 < 0$) decreasing exponentially with depth.



2. a) Relaxing the $kH \gg 1$ solution and again considering small trajectories $\vec{s} = (s, \xi)$, $k|\vec{s}| \ll 1$, we can write

$$\int \vec{s} = \int \vec{u}(\vec{x}_0) dt = \int \nabla \phi(\vec{x}_0) dt, \text{ giving}$$

$$\xi = -a \frac{\cosh k(z_0+H)}{\sinh kH} \sin(kx_0 - \omega t)$$


$$s = a \frac{\sinh k(z_0+H)}{\sinh kH} \cos(kx_0 - \omega t)$$

Solving the first equation for $\sin(kx_0 - \omega t)$ and the second for $\cos(kx_0 - \omega t)$ allows us to write

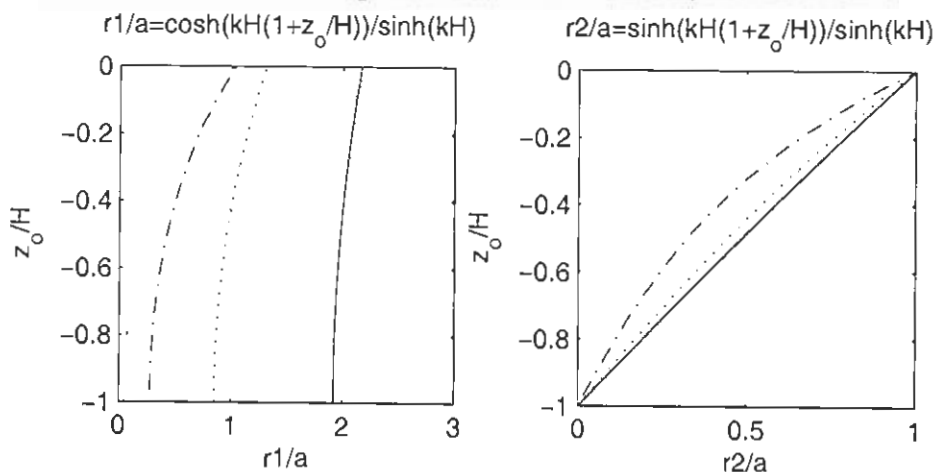
$$1 = \cos^2(kx_0 - \omega t) + \sin^2(kx_0 - \omega t) = \frac{\xi^2}{r_1^2} + \frac{s^2}{r_2^2} \quad \text{with}$$

$$r_1 = a \frac{\cosh k(z_0+H)}{\sinh kH}, \quad r_2 = a \frac{\sinh k(z_0+H)}{\sinh kH}$$

This is the equation for an ellipse, $\left(\frac{x}{r_1}\right)^2 + \left(\frac{y}{r_2}\right)^2 = 1$.

b) The two axes of the ellipse are r_1 and r_2 , defined above. 

c) Below r_1 and r_2 are plotted for several ratios of the horizontal wavelength to the depth ($\sim kH$).



d) Near the bottom, $r_2 \rightarrow 0$ while $r_1 \neq 0$, so the trajectories become horizontal oscillations. \longleftrightarrow

3. As shown in class, when $H \gg \lambda$ ($kH \gg 1$), $\omega = \sqrt{gk}$, and $C_g = \frac{d\omega}{dk} = \frac{1}{2} \sqrt{\frac{g}{k}} = \frac{1}{2} C_p$ [$C_p = \frac{\omega}{k}$], i.e., the group speed (speed of packet) is half the phase speed (speed of crests).